10/9/2025 Day 19. Solutions for time-independent Time-vidependent Hamiltonian. Let us specialize la lon care where the interactions are not dep. Since the Hamiltonian is Hermitean, ets eigenstells would provide a convenient representation for the state. KEY = EE - (10) or E = <E |Ĥ|E> - (11) (1E> E corresponds to energies Q llu sydem when in the date So we can write I Ich = I a CHIE> where aft)= \EII(+)>

(13)

Mariny (3) in (7) we get.

$$\sum_{E'} i \star \partial a_{e}^{(k)} |_{E'} = \sum_{E'} \hat{H}|_{E'}^{(k)} \alpha_{e}^{(k)} \\
= \sum_{E'} E'|_{E'}^{(k)} \alpha_{e}^{(k)}$$

$$= \sum_{E'} E'|_{E'}^{(k)}$$

at

$$\begin{array}{l}
\Rightarrow \quad a_{E}(t) = a_{E}(0) \exp(-i\frac{E}{E}t) \\
\downarrow^{1st} \text{ order linear homogeneous diffurities} \\
aquatin with initial condition }
a_{E}(t=0) = a_{E}(0)
\end{array}$$

$$\begin{array}{l}
a_{E}(t=0) = a_{E}(0) \\
\downarrow^{0}
\end{array}$$

$$\begin{array}{l}
\downarrow^{0}
\downarrow^{0}
\downarrow^{0}
\end{matrix}$$

$$\begin{array}{l}
\downarrow^{0}
\downarrow^{0}
\downarrow^{0}
\end{matrix}$$

Therefore, the standard protocol to solve the TDSE for any system is as Jollows:

- 1. Write down the Hamiltonian for the system H. (time-indep.)
- 2. Set up and solve the eigenvalue equation for the Hauvitonian (the teme-independent SE or TISE).

3. Determine the initial value of the energy coefficients given $|\Psi(0)\rangle$ (19) $Q_{E}(0) = \langle E|\Psi(0)\rangle$ (19)

4 Construct the time-dependent state as $|\Psi(t)\rangle = \sum_{E} Q_{E}(0) e^{-\frac{iEt}{\hbar}|E\rangle} - 20$

We will, for the most of this course, Jollow Mis procedure.

(20) says that any system is, at any instant, in a superposition of its energy eigenetales. This brings special importance to the TISE and, hence, we will be occupied with solvery this equation for many model systems.

However, (18) is still an abstract equality det us choose a representation to some et. The conventional choice is the position representation.

Let us, for the present, focus on as 1-particle system with a particle of mass on moving along the z-axis.

$$\hat{\mathcal{H}} = \frac{\hat{P}_{z}^{2}}{2m} + V(\hat{z}) - (\hat{z})$$

$$\hat{\mathcal{H}}(E) = E(E)$$

$$\hat{\mathcal{H}}|E\rangle = E|E\rangle$$

$$\Rightarrow \langle z|\hat{\mathcal{H}}|E\rangle = E\langle z|E\rangle - 22\rangle$$

$$LHS 9 22$$

$$= |\langle z|\hat{p}_{z}^{2}|E\rangle + \langle z|V(2)|E\rangle$$

$$= \sum_{z=1}^{2} \langle z|\hat{p}_{z}^{2}|E\rangle + \langle z|V(2)|E\rangle$$

$$= -\frac{\hbar^2}{2m} \frac{d^2}{dz^2} \langle xlE \rangle + V(z) \langle zlE \rangle$$

$$= \left(-\frac{\hbar^2}{2m} \frac{d^2}{dz^2} + V(z) \right) \psi(z)$$

$$= \left(-\frac{\hbar^2}{2m} \frac{d^2}{dz^2} + V(z) \right) \psi(z)$$

$$\begin{aligned}
&= \left[-\frac{1}{2m} \frac{d}{dx^{2}} + V(2) \right] \psi(2) \\
&= \left[-\frac{1}{2m} \frac{d}{dx^{2}} + V(2) \right] \psi(2) \\
&= \left[-\frac{1}{2m} \frac{d}{dx^{2}} + V(2) \right] \psi(2) \\
&= \left[-\frac{1}{2m} \frac{d}{dx^{2}} + V(2) \right] \psi(2) \\
&= \left[-\frac{1}{2m} \frac{d}{dx^{2}} + \frac{1}{2m} \frac{d}{dx^{2}} \right] \\
&= \left[-\frac{1}{2m} \frac{d}{dx^{2}} + \frac{1}{2m} \frac{d}{dx^{2}} \right] \\
&= \left[-\frac{1}{2m} \frac{d}{dx^{2}} + \frac{1}{2m} \frac{d}{dx^{2}} \right] \\
&= \left[-\frac{1}{2m} \frac{d}{dx^{2}} + \frac{1}{2m} \frac{d}{dx^{2}} \right] \\
&= \left[-\frac{1}{2m} \frac{d}{dx^{2}} + \frac{1}{2m} \frac{d}{dx^{2}} \right] \\
&= \left[-\frac{1}{2m} \frac{d}{dx^{2}} + \frac{1}{2m} \frac{d}{dx^{2}} \right] \\
&= \left[-\frac{1}{2m} \frac{d}{dx^{2}} + \frac{1}{2m} \frac{d}{dx^{2}} \right] \\
&= \left[-\frac{1}{2m} \frac{d}{dx^{2}} + \frac{1}{2m} \frac{d}{dx^{2}} \right] \\
&= \left[-\frac{1}{2m} \frac{d}{dx^{2}} + \frac{1}{2m} \frac{d}{dx^{2}} \right] \\
&= \left[-\frac{1}{2m} \frac{d}{dx^{2}} + \frac{1}{2m} \frac{d}{dx^{2}} \right] \\
&= \left[-\frac{1}{2m} \frac{d}{dx^{2}} + \frac{1}{2m} \frac{d}{dx^{2}} \right] \\
&= \left[-\frac{1}{2m} \frac{d}{dx^{2}} + \frac{1}{2m} \frac{d}{dx^{2}} \right] \\
&= \left[-\frac{1}{2m} \frac{d}{dx^{2}} + \frac{1}{2m} \frac{d}{dx^{2}} \right] \\
&= \left[-\frac{1}{2m} \frac{d}{dx^{2}} + \frac{1}{2m} \frac{d}{dx^{2}} \right] \\
&= \left[-\frac{1}{2m} \frac{d}{dx^{2}} + \frac{1}{2m} \frac{d}{dx^{2}} \right] \\
&= \left[-\frac{1}{2m} \frac{d}{dx^{2}} + \frac{1}{2m} \frac{d}{dx^{2}} \right] \\
&= \left[-\frac{1}{2m} \frac{d}{dx^{2}} + \frac{1}{2m} \frac{d}{dx^{2}} \right] \\
&= \left[-\frac{1}{2m} \frac{d}{dx^{2}} + \frac{1}{2m} \frac{d}{dx^{2}} \right] \\
&= \left[-\frac{1}{2m} \frac{d}{dx^{2}} + \frac{1}{2m} \frac{d}{dx^{2}} \right] \\
&= \left[-\frac{1}{2m} \frac{d}{dx^{2}} + \frac{1}{2m} \frac{d}{dx^{2}} \right] \\
&= \left[-\frac{1}{2m} \frac{d}{dx^{2}} + \frac{1}{2m} \frac{d}{dx^{2}} \right] \\
&= \left[-\frac{1}{2m} \frac{d}{dx^{2}} + \frac{1}{2m} \frac{d}{dx^{2}} \right] \\
&= \left[-\frac{1}{2m} \frac{d}{dx^{2}} + \frac{1}{2m} \frac{d}{dx^{2}} \right] \\
&= \left[-\frac{1}{2m} \frac{d}{dx^{2}} + \frac{1}{2m} \frac{d}{dx^{2}} \right] \\
&= \left[-\frac{1}{2m} \frac{d}{dx^{2}} + \frac{1}{2m} \frac{d}{dx^{2}} \right] \\
&= \left[-\frac{1}{2m} \frac{d}{dx^{2}} + \frac{1}{2m} \frac{d}{dx^{2}} \right] \\
&= \left[-\frac{1}{2m} \frac{d}{dx^{2}} + \frac{1}{2m} \frac{d}{dx^{2}} \right] \\
&= \left[-\frac{1}{2m} \frac{d}{dx^{2}} + \frac{1}{2m} \frac{d}{dx^{2}} \right] \\
&= \left[-\frac{1}{2m} \frac{d}{dx^{2}} + \frac{1}{2m} \frac{d}{dx^{2}} \right] \\
&= \left[-\frac{1}{2m} \frac{d}{dx^{2}} + \frac{1}{2m} \frac{d}{dx^{2}} \right] \\
&= \left[-\frac{1}{2m} \frac{d}{dx^{2}} + \frac{1}{2m} \frac{d}{dx^{2}} \right] \\
&= \left[-\frac{1}{2m} \frac{d}{dx^{2}} + \frac{1}{2m} \frac{d}{dx^{2}} \right] \\
&$$