

# Density Functional Theory - Levy-Lied Constrained Search Formulation

*Lecture 28*

---

CHM 652 / PHY 626

Electronic Structure of Materials

Varadharajan Srinivasan  
Dept. Of Chemistry  
IISER Bhopal

---



---

# Lecture Plan

---

- $N$  and  $v$  representable densities
- Levy-Lieb constrained search formulation of DFT



---

# $v$ representable densities

---

In deriving the Hohenberg-Kohn theorems we implicitly assumed that the densities of interest can be obtained by applying some external potential to an  $N$ -electron system. This was crucial in deriving the universal functional as well as the variational principle.

Densities associated with antisymmetric  $N$ -electron ground state wavefunctions of a Hamiltonian that includes an external potential are termed as  *$v$ -representable densities*.

Since  $F_{HK}$  is derived from such a wave function, it is also a functional of  $v$ -representable densities. That is, the variational principle is only valid for such densities.

Not all valid densities are necessarily  *$v$ -representable*.



---

# N representable densities

---

Densities that are derived from any antisymmetric  $N$ -electron wavefunction are termed as *N-representable densities*.

Mathematically such densities satisfy (see Parr and Yang)

$$n(\mathbf{r}) \geq 0, \quad \int n(\mathbf{r}) d^3r = N, \quad \int \left| \nabla n(\mathbf{r})^{\frac{1}{2}} \right|^2 d^3r < \infty$$

This is a weaker condition.

*Pure-state representable* : densities obtained from a ground state antisymmetric wavefunction (same as *v-representable*).

*Non-interacting pure-state representable* : As above except that the Hamiltonian has no interaction among the electrons.

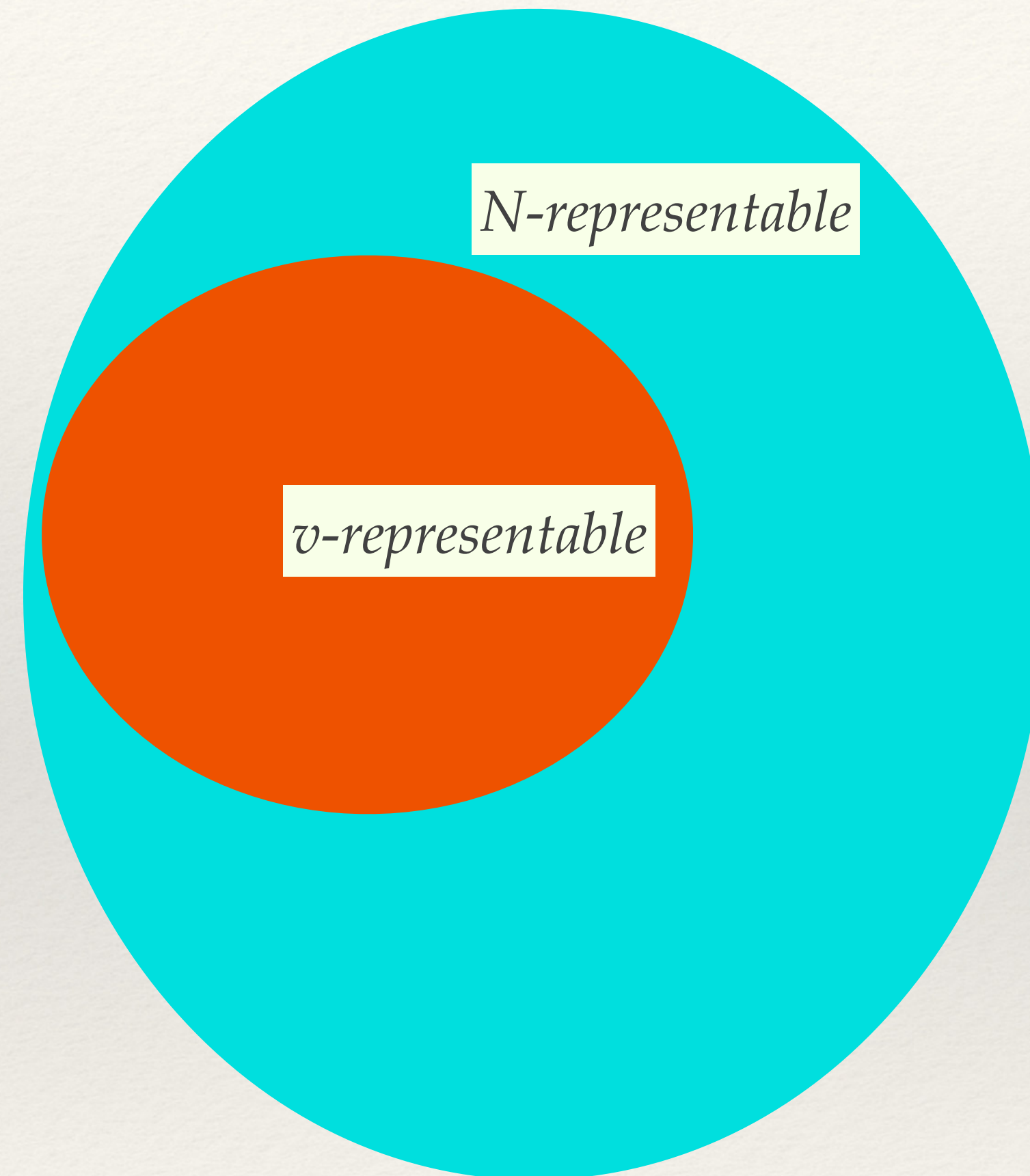
*Mixed-state representable* : densities obtained by mixing *pure-state* densities.



---

# $N$ and $v$ representable wavefunctions

---





# Levy constrained search formulation

The variation principle suggests a search over the *N-representable* space

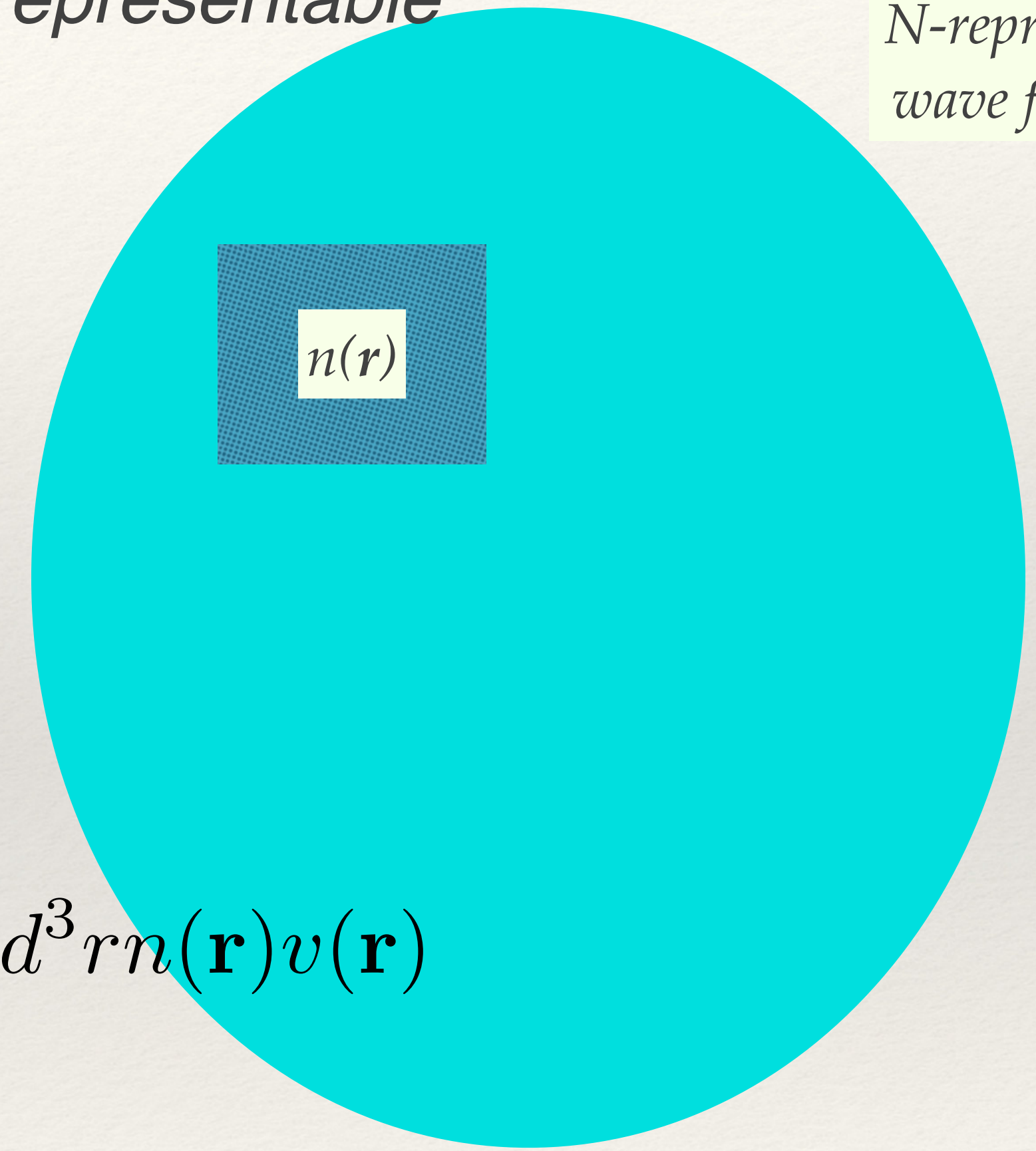
$$E_0 = \min_{\Psi} \langle \Psi | \hat{H} | \Psi \rangle$$

Suppose, we break this search into two steps.

1) Search for the minimum over a wave functions yielding a given density  $n(\mathbf{r})$

$$\begin{aligned} E[n] &= \min_{\Psi \rightarrow n(\mathbf{r})} \langle \Psi | \hat{H} | \Psi \rangle \\ &= \min_{\Psi \rightarrow n(\mathbf{r})} \langle \Psi | \hat{T} + \hat{W} | \Psi \rangle + \int d^3r n(\mathbf{r}) v(\mathbf{r}) \\ &= F_{LL}[n] + \int d^3r n(\mathbf{r}) v(\mathbf{r}) \end{aligned}$$

*N-representable  
wave functions*





# Levy constrained search formulation

2) Search for the minimum over all densities of the energy functional

$$E_0[n_0] = \min_{n(\mathbf{r})} \left( F_{LL}[n] + \int d^3r n(\mathbf{r}) \right)$$

Where  $n_0(\mathbf{r})$  is the minimising ground-state density.

Thus, the constrained search formulation extends the validity of the HK theorems to *N-representable* densities.

